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RESONANCE FUNCTIONS OF TWO-BODY SCHRÖDINGER OPERATORS

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We consider the Schrödinger operator $-\Delta + V$ in $L^2(\mathbb{R}^n)$, $n \ge 3$, where V is a short-range, dilation-analytic potential in an angle S_{α} . A resonance λ_0 appears as a discrete eigenvalue of the complex-dilated Hamiltonian [2], a pole of the S-matrix [3] and as a pole of the analytically continued resolvent, acting from an exponentially weighted space to its dual [4,5]. In [2] resonance functions are obtained as square-integrable eigenfunctions of the complex-dilated Hamiltonian, corresponding to the eigenvalue λ_0 , in [5] they are defined as certain exponentially growing solutions f of the Schrödinger equation $(-\Delta + V - \lambda_0)f = 0$. In [6] it is proved that for a dilation-analytic multiplicative potential V with resonance λ_0 , the resonance functions of [2] and [5] are simply the restrictions of one analytic, $L^2(s^{n-1})$ -valued function f on S_{α} to rays $e^{i\phi}\mathbb{R}^+$ with $2\phi > -\operatorname{Arg}\lambda_0$ and to \mathbb{R}^+ , respectively.

Moreover, the precise asymptotic behaviour $f(z) = e^{1K_0 z}$ $\frac{1-n}{z}(\tau + 0(|z|^{-\varepsilon}))$ with $\tau \in L^2(S^{n-1})$, where $k_0^2 = \lambda_0$, is established together with asymptotics for f'(z). These imply exponential decay in time of resonance states, defined as suitably cut-off resonance functions, as proved in [8].

In this note we shall give a brief account of results on resonance functions, referring for details to [5] and [6].

XIV-1

We introduce the weighted L^2 -spaces $L^2_{\delta,b} = L^2_{\delta,b}(\mathbb{IR}^n)$ for $\delta, b \in \mathbb{IR}$ by

$$L_{\delta,b}^{2} = \{f \mid \|f\| \|_{\delta,b}^{2} = \int_{IR^{n}} |f(x)|^{2} (1+r^{2})^{\delta} e^{2br} dx < \infty \}$$

where $x \in \operatorname{IR}^n$, r = |x|. The weighted Sobolev spaces are defined by

$$H^{2}_{\delta,b} = \{f \mid \|f\|^{2}_{2,\delta,b} = \sum_{|\alpha| \leq 2} \|D^{\alpha}f\|^{2}_{\delta,b} < \infty \}$$

We set $L_{\delta,0}^2 = L_{\delta}^2$, $H_{\delta,0}^2 = H^2$ and $h = L^2(S^{n-1})$, $S^{n-1} = {x \in IR^n \mid |x| = 1}$. We assume that the dimension $n \ge 3$

$$\mathbf{C}^+ = \{\mathbf{k} \in \mathbf{C} \mid \text{Im } \mathbf{k} > 0\}$$
, $\widetilde{\mathbf{C}}^+ = \overline{\mathbf{C}^+} \setminus \{0\}$.

 $B(H_1, H_2)$ and $C(H_1, H_2)$ denote the spaces of bounded and compact operators from H_1 into H_2 , respectively.

The free Hamiltonian H_0 in L^2 is defined for $u \in \mathcal{D}_{H_0} = H^2$ by $H_0 u = -\Delta u$ with resolvent $R_0(k) = (H_0 - k^2)^{-1} \in \mathcal{B}(L^2)$ for $k \in \mathbf{C}^+$.

The interaction Q is assumed to be a symmetric, short-range, S_{α} -dilation-analytic operator in L^2 . Thus, $Q \in C(H^2_{-\delta_0}, L^2_{\delta_0})$ for some $\delta_0 > \frac{1}{2}$, and if $\{U(\rho)\}$ is the dilation group on L^2 $\rho \in IR^+$ defined by n_{α}

$$(\mathbf{U}(\rho)\mathbf{f})(\mathbf{x}) = \rho^{\mathbf{\hat{2}}}\mathbf{f}(\rho\mathbf{x})$$

then the function $Q(\rho) = U(\rho)QU(\rho^{-1})$ on \mathbb{R}^+ has an analytic, $C(H^2_{-\delta_0}, L^2_{\delta_0})$ -valued analytic extension to the angle

$$\mathbf{S}_{\alpha} = \{ \rho \mathbf{e}^{\mathbf{i} \boldsymbol{\varphi}} \mid \rho > 0 , |\boldsymbol{\varphi}| < \alpha \}$$

Moreover, $Q(z) \in C(H^2_{-\delta_0,b}, L^2_{\delta_0,b})$ for all $b \in \mathbb{R}$. (This follows from $Q \in C(H^2_{-\delta_0}, L^2_{\delta_0})$ if Q is local).

The Hamiltonian $H = H_0 + Q$ is self-adjoint on $\mathcal{D}_H = H^2$, and associated with H is a self-adjoint, analytic family of type A, H(z), given by

$$H(z) = z^{-2} H_0 + Q(z)$$

and $H(\rho e^{i\phi}) = U(\rho)H(e^{i\phi})U(\rho^{-1})$, so $\sigma(H(z)) = \sigma(H(e^{i\phi}))$ for $\rho > 0$, $z = \rho e^{i\phi}$.

We define the operators H_z and their resolvents $R_z(k)$ by

$$H_z = H_0 + z^2 Q(z) = z^2 H(z)$$
, $R_z(k) = (H_z - k^2)^{-1}$

We note that $R_{z}(zk) = z^{-2}(H(z) - k^{2})^{-1}$.

We have $\sigma_e(H(z)) = e^{-2i\varphi} \overline{IR}^+$ and $\sigma_d(H(z)) \setminus IR \subset \{\lambda \mid -2\varphi < \operatorname{Arg} \lambda < 0\}$.

We define $R(\varphi)$ by $R(\varphi) = \{k \mid 0 > \operatorname{Arg} k > -\varphi, k^2 \in \sigma_d(H(z))\}$, $R = \bigcup R(\varphi)$. The points $\lambda = k^2$, $k \in R$, are called resonances. $0 < \varphi < \alpha$ For our analysis we need the following result, proved in [5]:

Lemma 1.1. For $\delta > 0$ let $S_{\alpha}^{\delta} = \{k \in S_{\alpha} \mid \text{Im}(e^{i(\alpha - \delta)}k) < \epsilon\}$. There exists S_{α} -dilation-analytic interactions V_{ϵ} and W_{ϵ} with $Q = V_{\epsilon} + W_{\epsilon}$, such that $H_{0} + V_{\epsilon}$ has no resonances outside $(S_{\alpha}^{\delta})^{2}$ and W_{ϵ} decays faster than any exponential. This holds with $W_{\epsilon} = g_{\epsilon} Qg_{\epsilon}$, where $g_{\epsilon}(r) = \exp(-\epsilon r^{\beta})$, $\beta = \frac{\pi}{2\kappa}$, for ϵ small.

Using Lemma 1.1 one can prove all results for fixed $\delta > 0$ with S_{α} replaced by $S_{\alpha} - S_{\alpha}^{\delta}$, using the splitting $Q = V_{\epsilon} + W_{\epsilon}$, and then let $\delta \neq 0$. To simplify the presentation, we assume from the outset (although this can strictly speaking not be obtained) that $H_1 = H_0 + V$ has no resonances and fix ϵ , setting $g = g_{\epsilon}$, W = Qg, V = Q - gW. We denote by H_{1z} , $R_{1z}(k)$ etc. the operators obtained by replacing Q by V.

XIV-3

Basic to our approach is an extended limiting absorption principle proved in [7] and generalized in [5] to non-symmetric, short-range potentials like Q_z . The idea is to consider $-\Delta$ and $-\Delta + Q_z$ as operators H_0^{-b} and H_z^{-b} acting in the space $L_{0,-b}^2$, $b \ge 0$. The spectrum of H_0^{-b} and the essential spectrum of H_z^{-b} coincide with the parabolic region $P_b = \{k^2 \mid |\text{Im}k| \le b\}$, and it is then proved that the resolvents $(H_0^{-b} - (a + ib + i\epsilon)^2)^{-1}$ and $(H_z^{-b} - (a + ib + i\epsilon)^2)^{-1}$ have boundary values as $\epsilon \neq 0$ in $\mathcal{B}(L_{\delta,-b}^2, H_{-\delta-b}^2)$ for $\frac{1}{2} < \delta \le \delta_0$, except at the so-called singular points.

The singular sets \sum_{z}^{c} , \sum_{z}^{r} and \sum_{z} are defined for $z = \rho e^{i\phi}$, $\phi > 0$, by

$$\sum_{z}^{c} = \{ \mathbf{k} \in \mathbf{C}^{+} \mid \mathbf{k}^{2} = \mathbf{z}^{2} \lambda , \lambda \in \sigma_{d}(\mathbf{H}(\mathbf{z})) \} ,$$
$$\sum_{z}^{r} = \mathbf{z} R \cap \mathbf{R}^{+} , \sum_{z} = \sum_{z}^{c} \cup \sum_{z}^{r} ,$$

and for $\varphi < 0$ by $\sum_{z}^{c} = -\overline{\sum_{z}^{c}}$ and similar for \sum_{z}^{r} and \sum_{z}^{r} . For $\varphi = 0$, $\sum_{z}^{c} = \sum_{z}^{c} \cup \sum_{z}^{r} = \{k \in \widetilde{\mathbf{C}}^{+} \mid k^{2} \in \sigma_{p}(\mathbf{H})\}$.

The extended limiting absorption principle for H_z can then be formulated as follows:

<u>Theorem 1.2</u>. For fixed $z \in S_{\alpha}$, $0 < \delta \leq \delta_{0}$, there exists a meromorphic $B(L^{2}_{\delta}, H^{2}_{-\delta})$ -valued function $R_{z}^{-}(k)$ in \mathbb{C}^{+} , continuous in $\mathbb{C}^{+} \setminus \Sigma_{z}$, such that for $k \in \mathrm{IR} \setminus \Sigma_{z}^{r} \cup \{0\}$

$$R_z(k) = e^{ikr}R_z(k+i0)e^{-ikr}$$

where

$$R_{z}(k + i0) = \lim_{\epsilon \neq 0} R_{z}(k + i\epsilon)$$

in the operator-norm topology of $\ B(L^2_{\ \delta},H^2_{-\delta})$, locally uniformly in k .

For $f \in L^2_{\delta,-b}$, $u = e^{-ikr} R_z^-(k) e^{ikr} f$ is the unique solution in $L^2_{\delta,-b}$ of the equation $(H_z^{-b} - k^2)u = f$, such that $\mathcal{D}u \in L^2_{\delta-1,-b}$, where b = Imk and

$$\mathcal{D}u = (\mathcal{D}_1 u, \dots, \mathcal{D}_n u)$$
, $\mathcal{D}_j = \frac{\partial}{\partial x_j} + \frac{n-1}{2r^2} x_j - ik \frac{x_j}{r}$

(the radiation condition).

<u>Proof</u>. We refer to [5] for the proof of the Theorem. It utilizes the result of [7] for H_0 , analytic Fredholm theory and control of the singular points using analyticity in k and z.

The trace operators $T_0^{}(k)$, $T_z^{}(k)\in B(L^2_{~\delta},h)~$ are defined for $z\in S_\alpha^{}$, by

$$(\mathbf{T}_{0}(\pm \mathbf{k})\mathbf{f})(\mathbf{k},\cdot) = (\mathcal{F}_{\pm}\mathbf{f})(\mathbf{k},\cdot) , \quad \mathbf{k} \in \mathbf{R}^{+}$$

n

where

$$(F_{\pm}f)(k,\omega) = (2\pi)^{-\frac{\pi}{2}} \int_{\mathbb{R}^{n}} e^{\pm ik\omega \cdot x} f(x) dx ,$$
$$T_{z}(k) = T_{0}(k) (1 - Q_{z}R_{z}(k+i0)) , \quad k \in \mathbb{R} \setminus \sum_{z}^{r}$$

We set

$$T_0(k) = T_0(k)e^{ikr}$$
, $T_2^+(k) = T_2(k)e^{ikr}$

The following result is proved in [5].

<u>Theorem 1.3</u>. For $\frac{1}{2} < \delta \leq \delta_0$, $z \in S_\alpha$, the $\mathcal{B}(L^2_{\delta}, h)$ -valued function $T_z^+(k)$ has a continuous extension to $\widetilde{\mathbb{C}}^+ \setminus \sum_z$ meromorphic in \mathbb{C}^+ with poles at $\sum_z^{\mathbf{C}}$. The function $T_{\overline{z}}^{+*}(\overline{k})$ defined for $k \in \widetilde{\mathbb{C}}^- \setminus (-\sum_z)$ is analytic in $\mathbb{C}^- \setminus (-\sum_z^{\mathbf{C}})$ and continuous in $\widetilde{\mathbb{C}}^- \setminus (-\sum_z)$ as a $\mathcal{B}(h, H^2_{-\delta})$ -valued function.

We recall the following formulas from the stationary scattering theory:

$$R(k + i0) = R(-k + i0) + \pi i k^{n-2} T^{*}(k)T(k) , k \in \mathbb{R}^{+} \setminus \sum_{r}$$
(1.2)

$$T(k) = S(k)RT(-k)$$
 (1.3)

where $(R\tau)(\omega) = \tau(\omega)$ for $\tau \in h$.

Inserting (1.3) in (1.2), we obtain

$$R(k + i0) = R(-k + i0) + \pi i k^{n-2} T^{*}(k) S(k) RT(-k)$$
(1.4)

The S-matrix S(k) of
$$(H_0, H)$$
 is given for $k \in IR^+ \setminus \sum_r$ by
S(k) = 1 - $\pi i k^{n-2} T_0(k) (Q - QR(k + i0)Q) T_0^*(k)$ (1.5)

and the S-matrix $S_1(k)$ of (H_0, H_1) by (1.5) with Q and R replaced by V and R_1 .

The following result is proved in [3].

<u>Theorem 1.4</u>. The S-matrix S(k) has a meromorphic extension $\tilde{S}(k)$ from \mathbb{IR}^+ to S_{α} with poles at R. The S-matrix $S_1(k)$ has an analytic extension $\tilde{S}_1(k)$ from \mathbb{IR}^+ to S_{α} . Moreover, for k > 0, $0 < \varphi < \alpha$, $\tilde{S}_1(ke^{-i\varphi}) = S_{1,e^{i\varphi}}(k)$, where $S_{1,e^{i\varphi}}(k)$ is the 1, $e^{i\varphi}(k)$ is the 1, $e^{i\varphi}(k)$.

From (1.4) and Theorem 1.2 we obtain

<u>Theorem 1.5</u>. For any b > 0 the $B(L_{0,b}^2, H_{0,-b}^2)$ -valued function R(k) has a meromorphic continuation $\widetilde{R}(k)$ from \mathbb{C}^+ across IR^+ to $S_{\alpha,b} = \{k \in S_{\alpha} \mid -b < Imk < 0\}$, given by

$$\widetilde{R}(k) = R(-k) + \pi i k^{n-2} T^*(\overline{k}) \widetilde{S}(k) T(-k)$$
(1.6)

The $B(L_{0,b}^2, H_{0,-b}^2)$ -valued function $R_1(k)$ has an analytic continuation $\widetilde{R}_1(k)$ from $\mathbb{C}^+ \sum_{1c} \operatorname{across} \mathbb{R}^+$ to $S_{\alpha,b}$, given by (1.6) with R,T and S replaced by R_1, T_1 and S_1 .

The functions $\widetilde{R}(k)$ and $\widetilde{R}_1(k)$ are connected by the analytically continued symmetrized resolvent equation

$$\widetilde{\mathbf{R}}(\mathbf{k}) = \widetilde{\mathbf{R}}_{1}(\mathbf{k}) - \widetilde{\mathbf{R}}_{1}(\mathbf{k})g(1 + W\widetilde{\mathbf{R}}_{1}(\mathbf{k})g)^{-1} W\widetilde{\mathbf{R}}_{1}(\mathbf{k})$$
(1.7)

The following result is proved in [5]:

Theorem 1.6. $\widetilde{R}(k)$ and $\widetilde{S}(k)$ have the same poles and of the same order in $S_{\alpha,b}$.

2. Resonance functions

Let k_0^2 be a resonance, and fix $b > -Imk_0$. Then k_0 is a pole of $\widetilde{R}(k) \in \mathcal{B}(L_{0,b}^2, H_{0,-b}^2)$, defined in Theorem 1.5. Let C be a circle separating k_0 from other poles and let

$$P = -\frac{1}{2\pi i} \int_C \widetilde{R}_2(k) dk^2$$

be the residue of $\tilde{R}_2(k)$ at k_0 , $P \in B(L_{0,b}^2, H_{0,-b}^2)$ is of finite rank.

The space F of resonance functions associated with k_0 is defined by

$$F = \{f \in R_p \mid (-\Delta + Q - k_0^2) f = 0\}$$

The following result is proved in [5]:

<u>Theorem 2.1</u>. F is the isomorphic image of $N(\tilde{S}^{-1}(k_0))$ and of $N(1 + W\tilde{R}_1(k_0)g)$ via the following maps:

$$N(\widetilde{S}^{-1}(k_0)) \ni \tau \rightarrow T^*(\overline{k}_0)\tau = f \in F$$
$$N(1 + W\widetilde{R}_1(k_0)g) \ni \phi \rightarrow \widetilde{R}_1(k_0)g\phi = f \in F$$

<u>Remark</u>. From the representation $f = T^*(\overline{k}_0)\tau$ we conclude by Theorem 1.3 and the uniqueness part of Theorem 1.2 that $f \in H^2_{-\delta}, -b_0 \sim L^2_{\delta-1}, -b_0$ for every $\delta > \frac{1}{2}$ and $b_0 = -Imk_0$. A further analysis yields precise asymptotic estimates. We first establish the analyticity properties, using the second isomorphism.

Applying (1.4) to the operator H_{1z} at a point zk_0 with Arg $zk_0 = 0$ and noting that by Theorem 1.4, $S_{1z}(zk_0) = \tilde{S}_1(k_0)$ we obtain

$$R_{1z}(zk_{0} + i0) = R_{1z}(-zk_{0} + i0) + \pi i(zk_{0})^{n-2}$$

$$T_{1z}^{*}(\overline{zk}_{0}) \widetilde{S}_{1}(k_{0}) RT_{1z}(-zk_{0})$$
(1.7)

By Theorems 1.2 and 1.3 we obtain from (1.7)

 $\begin{array}{ccc} \underline{\text{Theorem 2.2}}, & \text{The } & B(L_{\delta}^{2}, H_{-\delta}^{2}) \text{-valued function } e^{-izk_{0}r} \\ e^{-izk_{0}r} & \text{has an analytic extension from } \{z \in zk_{0} \mid \text{IR}^{+}\} \\ \text{to } \{z \in S_{\alpha} \mid \text{Arg } zk_{0} < 0\} \text{, given by} \end{array}$

$$\begin{array}{ccc} -izk_{0}r & -izk_{0}r & -izk_{0}r & -izk_{0}r \\ e & \widetilde{R}_{1z}(zk_{0})e & = e & R_{1z}(-zk_{0})e & + \\ \pi i(zk_{0})^{n-2} & T_{1z}^{*}(\overline{zk}_{0})\widetilde{S}_{1}(k_{0})RT_{1z}(-zk_{0}) \end{array}$$
(1.8)

Recalling that $W_z = Q_z g(rz)$, where $g(rz) = \exp\{-\epsilon(rz)^{\beta}\}$ with $\beta > 1$, we obtain from Theorem 2.2

<u>Theorem 2.3</u>. The $C(L^2)$ -valued function $W_z R_{1z}(zk_0)g(rz)$ has an analytic continuation from $\{z \in S_\alpha \mid \operatorname{Arg} zk_0 > 0\}$ to $\{z \in S_\alpha \mid \operatorname{Arg} zk_0 \leq 0\}$, given by $W_z \widetilde{R}_{1z}(zk_0)g(rz)$.

By standard dilation-analytic arguments $\sigma(W_z \tilde{R}_1(zk_0)g(rz))$ is constant. Let C be a circle separating -1 from the rest of $\sigma(W_z \tilde{R}_1(zk_0)g(rz))$ and set

$$P(z) = -\frac{1}{2\pi i} \int_{C} (-\lambda + W_z \widetilde{R}_{1z}(zk_0)g(rz))^{-1} d\lambda .$$

Then P(z) is a dilation-analytic $B(L^2)$ -valued function of z, and P(z) is a projection on the finite-dimensional algebraic null space of $1 + W_z \tilde{R}_{1z} (zk_0) g(rz)$. Let $\phi \in N(1 + W \tilde{R}_1 (k_0) g(rz))$ and pick an S_α -dilation-analytic vector η in L^2 such that

 ϕ = P(1)n . Then $\phi(z)$ = P(z)n(z) $\in N(1 + W_z \widetilde{R}_{1z}(zk_0)g(rz)$, and $\phi(z)$ is dilation-analytic.

We now obtain, using the second isomorphism of Theorem 2.1,

<u>Theorem 2.4</u>. Let $f \in F$. Then there exists an S_{α} -dilation-analytic, $H^2_{-\delta}$ -valued function $\chi(z)$, such that $f = e^{ik_0 r} \chi(1)$ and for Arg z $k_0 > 0$

 $f(z) = e^{ik_0 zr} \chi(z) \in N(H(z) - k_0^2) .$ Moreover, $\chi(z) \notin L^2_{\delta-1}$ for all $z \in S_{\alpha}, \delta > \frac{1}{2} .$

<u>Proof</u>. Define f(z) by

$$f(z) = \begin{cases} izk_{0}r & -izk_{0}r \\ e & R_{1z}^{+}(zk_{0}) & e & g(rz)\phi(z) , Imzk_{0} > 0 \\ izk_{0}r & (-izk_{0}r & -izk_{0}r) & izk_{0}r \\ e & R_{1z}^{-}(zk_{0})e & e & g(rz)\phi(z) , \end{cases}$$

$$Imzk_{0} \leq 0$$

where $R_{1z}^{+}(zk_{0})$ is defined similarly to $R_{1z}^{-}(zk_{0})$, replacing -b by b and $e^{\pm iar}$ with $e^{\mp iar}$ in Theorem 1.2. Clearly, f(z) is continuous for $zk_{0} \in \mathbb{R}^{+}$. By Theorem 1.2 and 2.2, $\chi(z) = e^{-izk_{0}r}f(z)$ is an analytic $H_{-\delta}^{2}$ -valued function in S_{α} .

It follows from the uniqueness part of Theorem 1.2 that $\chi(z) \notin L^2_{\delta-1}$ for $\operatorname{Im} zk_0 < 0$. The fact that $\chi(z) \notin L^2_{\delta-1}$ for $\operatorname{Im} zk_0 \ge 0$ then follows by the next Lemma, proved in [6]:

Lemma 2.5. Let $\chi(z)$ be an S_a-dilation-analytic vector, and define $h(\varphi)$ for $\varphi \in (-\alpha, \alpha)$ by

$$h(\varphi) = \inf\{s \mid \chi(e^{i\varphi}) \in L^2_{-s}\}$$
.

Then either $h(\varphi) \equiv -\infty$ or $h(\varphi) > -\infty$ and h is convex in $(-\alpha, \alpha)$.

Using this Lemma together with a recent result of Agmon [1], giving the precise asymptotic behaviour of f(z) for $\operatorname{Arg} zk_0 > 0$, we finally obtain the desired asymptotic estimates of f and f'. We refer to [6] for the proof.

<u>Theorem 2.6</u>. Assume that Q is an S_{A} -dilation-analytic multiplicative potential such that $|Q(z)(x)| \leq C|x|^{-1-\varepsilon}$ for $z \in S_{\alpha}$ and $|x| \geq R$. Let $f \in F$. Then f is an analytic, *h*-valued function $f(z, \cdot)$ on S_{α} of the form

$$f(z,\cdot) = e^{ik_0 z} \frac{1-n}{z} g(z,\cdot)$$

where

 $g(z, \cdot) = \tau + 0(|z|^{-\varepsilon})$ $g'(z, \cdot) = 0(|z|^{-1-\varepsilon})$

uniformly in any smaller angle S'_{α} for some $\varepsilon > 0$. Moreover, $\tau \in N(\tilde{S}^{-1}(k_0))$ and $f = CT^*(\bar{k}_0)\tau$, $C = k_0^{\frac{n-1}{2}} (-i)^{\frac{1-n}{2}} (2\pi)^{\frac{1}{2}}$.

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