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ON THE POLES OF THE SCATTERING MATRIX

FOR TWO CONVEX OBSTACLES

by Mitsuru IKAWA

§1. Introduction.

Let $\mathcal O$ be a bounded open set in $\mathbb R^3$ with smooth boundary Γ . We set $\Omega = \mathbb R^3 - \overline{\mathcal O}$. Suppose that Ω is connected. Consider the following acoustic problem

(1.1)
$$\begin{cases} \Box u(x,t) = \frac{\partial^2 u}{\partial t^2} - \Delta u = 0 & \text{in } \Omega \times (-\infty, \infty) \\ u(x,t) = 0 & \text{on } \Gamma \times (-\infty, \infty) \end{cases}$$

where $\Delta = \int_{\delta=1}^{3} \partial^2 / \partial x_j^2$. Denote by $\mathcal{S}(z)$ the scattering matrix for this problem. About the defintion of the scattering matrix, see for example Lax and Phillips[7,page 9]. The result I like to talk about is the following

Theorem 1. Let \mathcal{O}_j , j=1,2, be open and strictly convex sets in \mathbb{R}^3 with smooth boundary Γ_j , that is, the Gaussian curvature of Γ_j is positive everywhere on Γ_j . Suppose that $\overline{\mathcal{O}_1} \cap \overline{\mathcal{O}_2} = \emptyset$. Then the scattering matrix $\mathfrak{L}(z)$ for

$$\mathcal{O} = \mathcal{O}_1 \cap \mathcal{O}_2$$

satisfies the following:

(1) There exist positive constant c_0 and c_1 such that $\mathcal{S}(z)$ is holomorphic in

{z; Im
$$z \le c_0 + c_1$$
} - $\bigvee_{j=-\infty}^{\infty} D_j$

where

$$D_{j} = \{z; |z - z_{j}| \le C(1+|j|)^{-1/2}\},\$$
 $z_{j} = ic_{0} + \frac{\pi}{d}j, \quad d = dis(\mathcal{O}_{1}, \mathcal{O}_{2}).$

- (2) For large |j|, every D_{j} contains exactly one pole of $\mathcal{S}(z)$.
- (3) Denoting the pole in D by ζ_j we have an asymptotic expansion

(1.2)
$$\zeta_{j} \sim z_{j} + \beta_{1} j^{-1} + \beta_{2} j^{-2} + \cdots$$
 for $|j| \rightarrow \infty$

where β_1 , β_2 , are complex constants determined by \mathcal{O} .

(4) In D_{j} $\mathcal{L}(z)$ is of the form

(1.3)
$$\mathcal{S}(z)f = \frac{1}{z-\zeta_{j}}(f,\psi_{j})m_{j} + \mathcal{H}_{j}(z)f$$

for all $f \in L^2(S^2)$, where m_j , $\psi_j \in L^2(S^2)$ such that $m_j \neq 0$, $\psi_j \neq 0$ and (\cdot, \cdot) stands for the scalar product in $L^2(S^2)$, and $\mathcal{H}_j(z)$ is an $\mathcal{L}(L^2(S^2), L^2(S^2))^{\frac{1}{2}}$ valued holomorphic function in D_j .

Concerning the existence of non-purely imaginary poles of $\mathcal{L}(z)$, Bardos, Guillot and Ralston[1] proved under the same assumption as ours the existence of an infinite number of the poles in

$$\{z; \text{ Im } z \leq \epsilon \log(1+|z|)\}$$

for any $\varepsilon>0$. This result is generalized by Petkov[11] and Petkov and Stojanov[12] to a case of many strictly convex obstacles. For non-strictly convex obstacles Ikawa[5] showed an example of two

¹⁾ We denote by $\mathcal{L}(E,F)$ the set of all linear bounded mappings from E into F.

convex obstacles whose scattering matrix has a sequence of the poles converging to the real axis. On the other hand Lebeau[9] considered the distribution of poles for one strictly convex obstacle.

§2. Reduction of the problem.

Consider a boundary value problem with a parameter $\mu \in \mathbb{C}$

(2.1)
$$\begin{cases} (\mu^2 - \Delta)u(x) = 0 & \text{in } \Omega \\ u(x) = g(x) & \text{on } \Gamma \end{cases}$$

for $g \in C^{\infty}(\Gamma)$. For $\text{Re}\mu > 0$ (2.1) has a solution u uniquely in $\bigcap_{m>0} H^m(\Omega)$. We denote the solution by $U(\mu)g$. Then $U(\mu)$ is holomorphic in $\text{Re}\mu > 0$ as an $\mathcal{Z}(C^{\infty}(\Gamma), C^{\infty}(\overline{\Omega}))$ -valued function. We shall show the following theorems on $U(\mu)$.

Theorem 2.1. (i) $U(\mu)$ is prolonged analytically as an $\mathcal{J}(C^{\infty}(\Gamma),C^{\infty}(\overline{\Omega}))$ -valued function into

$$\left\{ \operatorname{Re} \mu \geqslant -c_0-c_1 \right\} - \bigcup_{j=-\infty}^{\infty} \operatorname{iD}_j.$$

(ii) Set for k∈R

$$G_k = \{ \mu \in \mathbb{C}; |\mu + ik| \le c_0 + c_1, Re\mu \ge -c_0 - (\log(1+|k|))^{-1} \}.$$

Then for large |k|, U(μ) is represented in $G_k \cap \{Re\mu > 0\}$ as

(2.2)
$$U(\mu) = \frac{\beta(x,k,\mu)}{\Phi(\mu) - \gamma(k,\mu)} F(k,\mu) + V(k,\mu).$$

Here

(a) $\beta(\bullet,k,\mu)$ is $C^{\infty}(\overline{\Omega})$ -valued holomorphic function in G_k .

(b)
$$\mathcal{P}(\mu) = 1 - \lambda \lambda e^{-2d\mu}, \quad 0 < \lambda, \lambda < 1.$$

(c) For any N positive integer

$$| \gamma(k,\mu) - \sum_{1 \leq k \leq N} \sum_{0 \leq k \leq N} \gamma_{k,h} k^{-k} (\mu + ik)^{h} | \leq C_{N} |k|^{-N}$$

where $\gamma_{\text{l,h}}$ are complex constants.

- (d) $F(k,\mu)$ is an $\mathcal{L}(L^2(\Gamma),\mathbb{C})$ -valued holomorphic function in G_k .
- (f) $V(k,\mu)$ is an $\mathcal{Z}(C^{\infty}(\Gamma),C^{\infty}(\overline{\Omega}))$ -valued holomorphic function in G_k .

<u>Corollary.</u> $U(\mu)$ is prolonged analytically as $\mathcal{L}(C^{\infty}(\Gamma), C^{\infty}(\overline{\Omega}))$ valued function into

$$\bigcup_{\substack{|\mathbf{k}|: \text{large}}} (G_{\mathbf{k}} - \{\mu; \mathcal{P}(\mu) - \gamma(\mathbf{k}, \mu) = 0\}).$$

Theorem 2.2. Suppose that $\mu \in G_k$ and $\mathcal{P}(\mu) - \gamma(k,\mu) = 0$ for |k| large. Then we have

 $dim\{u; \mu\text{-outgoing solution of (2.1) for g=0}\} = 1.$

Note that the zeros of $\mathcal{P}(\mu)=0$ are $\{iz_j, j=0,\pm1,\pm2,\cdots\}$ and

$$\left|\frac{d}{d\mu}(\mathcal{P}(\mu) - \gamma(k,\mu))\right|_{\mu=iz_{j}} \geqslant d - C|k|^{-1}.$$

By setting $k=-\frac{\pi}{d}$ j we have that $\mathcal{P}(\mu)-\gamma(k,\mu)=0$ has only one zero in iD, and it is simple. Denote it by i ζ_j and we see that ζ_j has an asymptotic expansion (1.2).

Theorem 1 is immediately derived from Theorems 2.1 and 2.2 if we recall the relationships between $\mathcal{S}(z)$ and $U(\mu)$ shown in Lax and Phillips[7], especially Theorem 5.1 of Chapter V, which says that $\mathcal{S}(z)$ has a pole at exactly those points z such that $\mu=iz$ is a pole of $U(\mu)$.

§3. Sketch of the proofs of Theorems 2.1 and 2.2.

3.1. Asymptotic solutions for oscillatory boundary data.

Let $a_{j} \in \Gamma_{j}$ be the points verifying

$$|a_1 - a_2| = dis(O_1, O_2).$$

Denote by $S_{j}(\delta)$ for $\delta > 0$ a connected component containing a_{j} of $S_{j} \cap \{x; \ dis(x,L) = \delta\}$

where L is a straight line passing a_1 and a_2 , and denote by $\omega(\delta)$ a domain surrounded by $\{x; dis(x,L)=\delta\}$ and $S_j(\delta)$, j=1,2. Let $\upsilon_k(x)$ be a smooth function satisfying

$$v_{k}(x) = \begin{cases} 1 & \text{for } x \in S_{1}(k^{-\epsilon}) \\ 0 & \text{for } x \notin S_{1}((1+\delta)k^{-\epsilon}) \end{cases}$$

for some $\delta > 0$, $\epsilon > 0$ small constants. Let $h(t) \in C^{\infty}(0,d/2)$ satisfying $h(t) \ge 0$ and $\int h(t) dt = 1$. Set

(3.1)
$$m(x,t;k) = e^{ik(\psi(x)-t)} w(x)h(t-j(x))$$

where $\psi \in C^{\infty}(S_1(\delta_0))$ is a real valued function satisfying some conditions and j(x) a fixed smooth function determined by \mathcal{O} . We construct a sequence of functions of the form

(3.2)
$$u_{q}(x,t;k) = e^{ik(y_{q}(x)-t)N} \sum_{j=0}^{N} v_{j,q}(x,t;k)(ik)^{-j}$$

(I) ${m g}_{\rm q}$, q=0,1,... are determined successively by

$$\left\{ \begin{array}{ll} |\nabla \boldsymbol{\mathcal{G}}_0| = 1 & \text{in } \omega(\delta) \\ \\ \boldsymbol{\mathcal{G}}_0 = \psi & \text{and } \partial \boldsymbol{\mathcal{G}}_0/\partial n > 0 & \text{on } S_1(\delta), \end{array} \right.$$

$$\begin{cases} |\nabla \mathcal{G}_1| = 1 & \text{on } \omega(\delta) \\ \mathcal{G}_1 = \mathcal{G}_0 & \text{and } \partial \mathcal{G}_1/\partial n > 0 & \text{on } S_2(\delta), \\ |\nabla \mathcal{G}_2| = 1 & \text{on } \omega(\delta) \\ \mathcal{G}_2 = \mathcal{G}_1 & \text{and } \partial \mathcal{G}_2/\partial n > 0 & \text{on } S_1(\delta), \\ \vdots & \vdots & \vdots \end{cases}$$

(II) On amplitude functions.

Set

$$T_{q} = 2 \frac{\partial}{\partial t} + 2 \nabla \varphi_{q} \cdot \nabla + \Delta \varphi_{q}.$$

 $v_{0,q}$, q=0,1,2, · · · · are defined successively by

$$\begin{cases} T_0 v_{0,0} = 0 & \text{in } \omega(\delta) \times \mathbb{R} \\ v_{0,0} = f(x,t) & \text{on } \Gamma_1 \times \mathbb{R} \end{cases}$$

where f(x,t)=w(x)h(t-j(x)), and for $p \ge 1$

$$\begin{cases} T_{2p-1}v_{0,2p-1} = 0 & \text{in } \omega(\delta) \times \mathbb{R} \\ v_{0,2p-1} = v_{0,2p-2} & \text{on } \Gamma_{2} \times \mathbb{R}, \\ T_{2p}v_{0,2p} = 0 & \text{in } \omega(\delta) \times \mathbb{R} \\ v_{0,2p} = v_{k}(x)v_{0,2p-1} & \text{on } \Gamma_{1} \times \mathbb{R}. \end{cases}$$

Next for $j \ge 1$, $v_{j,q}$, $q=0,1,2,\cdots$ are defined successively for all $p \ge 0$ by

$$\begin{cases} T_{2p}v_{j,2p} = \square v_{j-1,2p} & \text{in } \omega(\delta) \times \mathbb{R} \\ v_{j,2p} = 0 & \text{on } \Gamma_1 \times \mathbb{R}, \end{cases}$$

$$\begin{cases} T_{2p+1}v_{j,2p+1} = \square v_{j-1,2p+1} & \text{in } \omega(\delta) \times \mathbb{R} \\ v_{j,2p+1} = v_{j,2p} & \text{on } \Gamma_2 \times \mathbb{R}. \end{cases}$$

On the asymptotic behavior of \mathcal{G}_q , $v_{j,q}$ for $q \longrightarrow \infty$, we have the following Lemmas.

Lemma 3.1. It holds that

$$| \mathcal{G}_{2p} - (\mathcal{G}_{\infty} + 2dp + d_0) |_{m} \le C_{m}^{\alpha^{2p}}$$

 $| \mathcal{G}_{2p+1} - (\mathcal{G}_{\infty} + (2p+1)d + d_0) |_{m} \le C_{m}^{\alpha^{2p}}$

where $oldsymbol{arphi}_{\infty}$, $oldsymbol{\widetilde{arphi}}_{\infty}$ are functions independent of ψ , and they verify

$$|\nabla \mathcal{G}_{\infty}| = 1$$
 in $\omega(\delta)$ and $\mathcal{G}(a_1) = 0$, $|\nabla \widetilde{\mathcal{G}}_{\infty}| = 1$ in $\omega(\delta)$ and $\widetilde{\mathcal{G}}(a_2) = 0$,

and d_0 is a constant depending on ψ , α is a positive constant < 1.

Lemma 3.2. It holds that

$$\begin{aligned} |\mathbf{v}_{\mathtt{j},2p}(\mathtt{x},\mathtt{t};\mathtt{k}) - \mathtt{bw}(\mathtt{A}) (\lambda \widetilde{\lambda})^{p} & \mathbf{v}_{\mathtt{j},\infty}(\mathtt{x},\mathtt{t-2pd-j}(\mathtt{A})-\mathtt{d}_{\infty};\mathtt{k})|_{\mathtt{m}} \\ & \leq C_{\mathtt{j},\mathtt{m}}(\alpha \lambda \widetilde{\lambda})^{p} \mathtt{M}_{\mathtt{m}+2\mathtt{j}}, \\ |\mathbf{v}_{\mathtt{j},2p+1}(\mathtt{x},\mathtt{t};\mathtt{k}) - \mathtt{bw}(\mathtt{A}) (\lambda \widetilde{\lambda})^{p} & \widetilde{\mathbf{v}}_{\mathtt{j},\infty}(\mathtt{x},\mathtt{t-2pd-j}(\mathtt{A})-\mathtt{d}_{\infty};\mathtt{k})|_{\mathtt{m}} \\ & \leq C_{\mathtt{j},\mathtt{m}}(\alpha \lambda \widetilde{\lambda})^{p} \mathtt{M}_{\mathtt{m}+2\mathtt{j}}, \end{aligned}$$

where λ , $\widetilde{\lambda}$ are constants determined by $\mathcal O$ such that $0 < \lambda$, $\widetilde{\lambda} < 1$,

$$M_{\ell} = k^{\epsilon \ell} \sum_{\beta \mid \beta \mid < \ell} \sup_{\Gamma_{x} \times \mathbf{R}} |D_{x,t}^{\beta}|,$$

 $v_{j,\infty}$ and $\widetilde{v}_{j,\infty}$ are functions of the form

$$v_{j,\infty}(x,t;k) = \sum_{k=0}^{2j} a_{j,k}(x,k)h^{(k)}(t-j(x)),$$

$$\widetilde{v}_{j,\infty}(x,t;k) = \sum_{k=0}^{2j} \widetilde{a}_{j,k}(x,k)h^{(k)}(t-\widetilde{j}(x)),$$

and b is a constant depending on ψ , A is a point in $S_1(\delta)$ depending on ψ .

Remark that we have

$$\Box u_{q} = e^{ik(\mathbf{g}_{q}-t)} (ik)^{-N} \Box_{V_{N,q}}.$$

Next we construct by a usual method asymptotic solutions for

(3.3)
$$\begin{cases} \Box u = 0 & \text{in } \omega \times \mathbb{R} \\ u = (1 - v_k(x)) u_{2p}(x, t; k) & \text{on } \Gamma_1 \times \mathbb{R} \end{cases}$$

Denote the asymptotic solution by u_{2p}^{\prime} . Extend $\Box (u_{2p}^{\prime} + u_{2p}^{\prime})$ and $\Box u_{2p+1}^{\prime}$ by a fixed manner into O so that these are smooth in $\mathbb{R}^3 \times \mathbb{R}$, and denote by $u_q^{\prime\prime}$ the solution of

where we set $u_{2p+1}^{\prime}=0$. By taking account of the continuity from u_q to u_q^{\prime} and $u_q^{\prime\prime}$ we have a Lemma of the type Lemma 3.2 on the convergence of u_q^{\prime} and $u_q^{\prime\prime}$. Set

$$r_q = u_q + u_q' + u_q''$$

and we have

Lemma 3.3. It holds that

$$\begin{split} \left| \mathbf{r}_{2p}(\mathbf{x},\mathsf{t};\mathbf{k}) - \mathbf{bw}(\mathbf{A}) \mathrm{e}^{-\mathrm{i}\mathbf{k}\,(\mathtt{j}\,(\mathbf{A}) + \mathrm{d}_{\infty})} \, \mathrm{e}^{\mathrm{i}\mathbf{k}\mathrm{d}_{0}} \, \left(\lambda\widetilde{\lambda}\right)^{p} \\ & \cdot \mathbf{r}_{\infty}(\mathbf{x},\mathsf{t} - 2\mathrm{p}\mathrm{d} - \mathrm{j}\,(\mathbf{A}) - \mathrm{d}_{\infty};\mathbf{k}) \, \big|_{m} \leqslant \, C_{m} \left(\alpha\lambda\widetilde{\lambda}\right)^{p} \, \mathbf{k}^{m+1} \\ \\ \left| \mathbf{r}_{2p+1}(\mathbf{x},\mathsf{t};\mathbf{k}) - \mathbf{bw}(\mathbf{A}) \mathrm{e}^{-\mathrm{i}\mathbf{k}\,(\mathtt{j}\,(\mathbf{A}) + \mathrm{d}_{\infty})} \, \, \mathrm{e}^{\mathrm{i}\mathbf{k}\mathrm{d}_{0}} \, \left(\lambda\widetilde{\lambda}\right)^{p} \\ & \cdot \widetilde{\mathbf{r}}_{\infty}(\mathbf{x},\mathsf{t} - 2\mathrm{p}\mathrm{d} - \mathrm{j}\,(\mathbf{A}) - \mathrm{d}_{\infty};\mathbf{k}) \, \big|_{m} \leqslant \, C_{m} \left(\alpha\lambda\widetilde{\lambda}\right)^{p} \, \mathbf{k}^{m+1}, \end{split}$$

where r_{∞} , \widetilde{r}_{∞} are functions independent of ψ and w.

Set

$$r(x,t;k) = \sum_{k=0}^{\infty} (-1)^{q} r_{q}(x,t;k).$$

Evidently it holds that

$$\Box r = 0$$
 in $\Omega \times \mathbb{R}$.

We consider the Laplace transformation of r in t, that is,

(3.4)
$$\hat{r}(x,\mu;k) = \int e^{-\mu t} r(x,t;k) dt.$$

We have from Lemma 3.3 the following

<u>Proposition 3.4.</u> Let $Re\mu > 0$. Then (3.4) converges and we have a representation of $\hat{\mathbf{r}}(\mathbf{x},\mu;\mathbf{k})$

(3.5)
$$\hat{r}(x,\mu;k)$$

$$= bw(A)e^{-(j(A)+d_{\infty})(\mu+ik)} e^{ikd_0} \mathcal{P}(\mu)^{-1} r_{\infty}(x,\mu;k) + s(x,\mu;k),$$
 where $\widehat{r}_{\infty}(x,\mu;k)$ is an entire function in μ independent of ψ and $\widehat{s}(x,\mu;k)$ is holomorphic in $\text{Re}\mu > -c_0-c_1$. Moreover we have

on Γ_1

$$\hat{r}(x,\mu;k) - e^{-(\mu+ik)j(x)} e^{ik\psi(x)} w(x) \hat{h}(\mu+ik)$$

$$= e^{ikd_0} bw(A)e^{-(j(A)+d_\infty)(\mu+ik)} \{ v_k(x) \sum_{k \in \mathbb{N}} \sum_{0 \le h \le N} a_{j,h}(x)(ik)^{-j} \}$$

$$x(\mu+ik)^{h} \hat{h}(\mu+ik) + a_{1}(x,k,\mu) \} \mathcal{P}(\mu)^{-1} + e_{1}(x,\mu;k),$$

and on Γ_2

$$\hat{r}(x,\mu;k) = e^{ikd_0} \frac{1}{p(\mu)} bw(A) e^{-(j(A)+d_\infty)(\mu+ik)} a_2(x,k,\mu) + e_2(x,k,\mu),$$

where $a_{j,h}(x)$ are smooth functions on $S_1(\delta)$, a_1 and a_2 are entire functions independent of ψ and w having an estimate

$$\sup_{\mathbf{x} \in \Gamma_{:}} \left| \mathbf{a}_{j}(\mathbf{x}, \mathbf{k}, \boldsymbol{\mu}) \right| \leq C_{N,R} \left| \mathbf{k} \right|^{-N}$$

and e_1 and e_2 are holomorphic in $Re\mu > -c_0-c_1$ and satisfy

$$\begin{aligned} \left| \mathbf{e}_{1}(\mathbf{x}, \mathbf{k}, \boldsymbol{\mu}) \right| & \leq \begin{cases} \left| \mathbf{k} \right|^{-1} & \text{on } \mathbf{S}_{1}((1+\delta) \left| \mathbf{k} \right|^{-1}) \\ \left| \mathbf{k} \right|^{-N} & \text{on } \Gamma_{1} \setminus \mathbf{S}_{1}((1+\delta) \left| \mathbf{k} \right|^{-1}), \end{aligned}$$

$$\left| \mathbf{e}_{2}(\mathbf{x}, \mathbf{k}, \boldsymbol{\mu}) \right| \leq \left| \mathbf{k} \right|^{-N} & \text{on } \Gamma_{2}.$$

3.2. Reduction to an integral equation on Γ_1 .

Suppose that Γ_1 is represented as $\mathbf{x}(\sigma) = (\sigma_1, \sigma_2, \mathbf{x}_3, (\sigma_1, \sigma_2))$ near \mathbf{a}_1 . Let $\mathbf{g}(\mathbf{x}) \in C_0^{\infty}(S_1(\delta_0))$. Then

$$g(\mathbf{x}(\sigma)) = (2\pi)^{-2} \int e^{\mathbf{i}\mathbf{k}\sigma \cdot \xi} \widehat{g}(\mathbf{k}\xi) k^{2} d\xi$$
$$= (2\pi)^{-2} \mathbf{w}(\mathbf{x}(\sigma)) \int e^{\mathbf{i}\mathbf{k}\sigma \cdot \xi} \widehat{g}(\mathbf{k}\xi) k^{2} d\xi$$

where $w(x) \in C_0^{\infty}(S_1(2\delta_0))$ such that w(x)=1 on $S_1(\delta_0)$, and $\hat{g}(\xi) = \int e^{-i\sigma \cdot \xi} g(x(\sigma)) d\sigma.$

If we define $\widetilde{U}_1(k,\mu)$ an operator from $L^2(S_1(\delta_0))$ into $C^{\infty}(\overline{\Omega})$ by

$$(\widetilde{\mathbb{U}}_{1}(k,\mu)g)(x) = (2\pi)^{-2} \int u(x,\xi;k,\mu)\widehat{g}(k\xi)k^{2}d\xi$$

where $u(x,\xi;k,\mu)$ denotes $\hat{r}(x,\mu,k)/\hat{h}(\mu+ik)$ constructed for $\psi(x(\sigma))$ = $\sigma \cdot \xi$. Then we have from Proposition 3.4

Proposition 3.5. $\tilde{U}_{1}(k,\mu)$ is of the form

(3.6)
$$\widetilde{U}_{1}(k,\mu)g = \frac{\widehat{r}_{\infty}(x,\mu;k)}{p(\mu)} F_{0}(k,\mu)g + S(k,\mu)g$$

where

(3.7)
$$F_{0}(k,\mu)g = (2\pi)^{-2} \int b(\xi)w(A(\xi))e^{-(j(A(\xi))+d_{\infty}(\xi))(\mu+ik)} \cdot e^{ikd_{0}(\xi)} \hat{g}(k\xi)k^{2}d\xi,$$

 $S(k,\mu)$ is $\mathcal{L}(L^2(S_1(\delta_0)),C^{\infty}(\overline{\Omega}))$ -valued holomorphic function in $Re\mu > -c_0-c_1$. Moreover it holds that

$$(3.8) (\mu^2 - \Delta)\widetilde{U}_1 g = 0 in \Omega,$$

(3.9)
$$\widetilde{U}_{1}g = g - \frac{\alpha(x,k,\mu)}{\mathcal{P}(\mu)} F_{0}(k,\mu)g - E(k,\mu)g \quad \text{on} \quad \Gamma_{1}$$

(3.10)
$$\widetilde{\mathbf{U}}_{1}g = \frac{\widetilde{\alpha}(\mathbf{x},\mathbf{k},\mu)}{\mathbf{\mathcal{P}}(\mu)} \mathbf{F}_{0}(\mathbf{k},\mu)g + \widetilde{\mathbf{E}}(\mathbf{k},\mu)g \quad \text{on} \quad \Gamma_{2},$$

(3.11)
$$|\alpha(x,k,\mu)F_0(k,\mu)g| \le C|k|^{-\epsilon} \|g\|_{L^2(\Gamma_1)}$$

(3.12)
$$\|E(k,\mu)g\|_{L^{2}(\Gamma_{1})} \leq C|k|^{-\epsilon} \|g\|_{L^{2}(\Gamma_{1})}$$

(3.13)
$$|\tilde{\alpha}(x,k,\mu)F_0(k,\mu)g| \leq C|k|^{-N} ||g||_{L^2(\Gamma_1)}$$

(3.14)
$$\|E(k,\mu)g\|_{L^{2}(\Gamma_{2})} \le C|k|^{-N} \|g\|_{L^{2}(\Gamma_{1})}$$

Note that the solution U2h of

$$\begin{cases} (\mu^2 - \Delta)u = 0 & \text{in } \mathbb{R}^3 - \overline{\mathfrak{O}}_2 \\ u = h & \text{on } \Gamma_2 \end{cases}$$

is continued into $\{\mu; Re\mu \geqslant -a \log(|\mu|+1)\}$ for some a > 0. Then

$$\mathbf{U}_{1}(\mathbf{k},\boldsymbol{\mu})\mathbf{g} = \widetilde{\mathbf{U}}_{1}(\mathbf{k},\boldsymbol{\mu})\mathbf{g} - \mathbf{U}_{2}(\boldsymbol{\mu}) (\widetilde{\mathbf{U}}_{1}(\mathbf{k},\boldsymbol{\mu})\mathbf{g}|_{\Gamma_{2}})$$

is also of the form (3.6) and satisfies (3.8), (3.9), (3.11) and (3.12), and

(3.10)'
$$U_1(k,\mu)g = 0$$
 on Γ_2 .

Remark. We can extend the definition of $U_1(k,\mu)$ for any f $\in L^2(\Gamma_1)$ by using the argument in §8 of [2]. Hereafter we denote by U_1 the extended one.

3.3. Representation of $U(\mu)$.

Lemma 3.6. Let H and E be linear operators with $\|H\|$, $\|E\| < 1/2$. Then we have

$$(I - H - E)^{-1} = I + \mathcal{C}_1 + \mathcal{C}_2$$

where

$$C_1 = \mathcal{H} + \mathcal{H}E + \mathcal{H}E\mathcal{H} + \mathcal{H}E\mathcal{H}E + \cdots,$$

$$C_2 = E + E\mathcal{H} + E\mathcal{H}E + E\mathcal{H}E\mathcal{H} + \cdots,$$

$$E = E + E^2 + E^3 + \cdots,$$

$$\mathcal{H} = H + H^2 + H^3 + \cdots.$$

Pose

$$H(k,\mu)g = \frac{\alpha(x,k,\mu)}{\mathcal{P}(\mu)} F_0(k,\mu)g.$$

An application of the above lemma gives

(3.15)
$$(I - H - E)^{-1} = (I + \mathcal{E}) + \frac{(I + \mathcal{E})\alpha}{\mathcal{P}(\mu) - \gamma} F_0(I + \mathcal{E})$$

where

$$\gamma\left(\mathbf{k},\boldsymbol{\mu}\right) \ = \ \mathbf{F}_{0}\left(\mathbf{k},\boldsymbol{\mu}\right)\left(\left(\mathbf{I} + \boldsymbol{\xi}(\mathbf{k},\boldsymbol{\mu})\right)\boldsymbol{\alpha}\left(\bullet,\mathbf{k},\boldsymbol{\mu}\right)\right).$$

Evidently we have in $Re\mu > 0$

(3.16)
$$U(\mu) = U_{\eta}(k,\mu) (I-H(k,\mu)-E(k,\mu))^{-1}.$$

Then a substitution of (3.15) into (3.16) gives

$$U(\mu) = \frac{r_{\infty}(x,k,\mu)}{p(\mu) - \gamma(k,\mu)} F_{0}(k,\mu) (I + \xi(k,\mu)) + S(k,\mu) (I + \xi(k,\mu)).$$

By posing

$$\mathbf{F}\left(\mathbf{k},\boldsymbol{\mu}\right) \ = \ \mathbf{F}_{0}\left(\mathbf{k},\boldsymbol{\mu}\right)\left(\mathbf{I} + \ \boldsymbol{\xi}(\mathbf{k},\boldsymbol{\mu})\right),$$

 $V(k,\mu) = S(k,\mu) (I + \xi(k,\mu)),$

we have a representation (2.2).

References

- [1] C.Bardos, J.C.Guillot and J.Ralston, La relation de Poisson pour l'équation des ondes dans un ouvert non borné. Application à la théorie de la diffusion, Comm. Partial Diff. Equ., 7(1982), 905-958.
- [2] M.Ikawa, Decay of solutions of the wave equation in the exterior of two convex obstacles, Osaka J.Math., 19(1982), 459-509.
- [3] ——, On the poles of the scattering matrix for two strictly convex obstacles, J.Math.Kyoto Univ., 23(1983), 127-194.
- [4] ——, On the poles of the scattering matrix for two strictly convex obstacles: Addendum, J.Math.Kyoto Univ., 23(1983),795-802.
- [5] ———, Trapping obstacles with a sequence of poles of the scattering matrix converging to the real axis, to appear in Osaka J.Math.
- [6] ———, Precise informations on the poles of the scattering matrix for two strictly convex obstacles, in preparation.
- [7] P.D.Lax and R.S.Phillips, Scattering theory, Academic Press, New York, (1967).
- [8] ______, A logarithmic bound on the location of the scattering matrix, Arch.Rat.Mech.and Anal., 40(1971), 268-280.
- [9] G.Lebeau, to appear.
- [10] R.Melrose, Polynomial bound on the distribution of poles in scattering by obstacles, Journées "Equations aux dérivées partielles", Soc.Math.France, (1984).

- [11] V.M.Petkov, Propriétés génériques des rayons réfléchissants et applications aux problèmes spectraux, Séminaire Bony-Sjöstrand-Meyer, 1984-1985, Exposé n°XII.
- [12] V.M.Petkov and L.Stojanov, Periods of multiple reflecting geodesics and inverse spectral problems, preprint.